Finite-Dimensional Reduction of Systems of Nonlinear Diffusion Equations

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Abstract—We present a class of one-dimensional systems of nonlinear parabolic equations for which the phase dynamics at large time can be described by an ODE with a Lipschitz vector field in \mathbb{R}^n . In the considered case of the Dirichlet boundary value problem, the sufficient conditions for a finite-dimensional reduction turn out to be much wider than the known conditions of this kind for a periodic situation.

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1. INTRODUCTION

One of the main problems in the study of evolution equations is related to the description of the final (long-time) behavior of their solutions. We consider systems of diffusion equations with the Dirichlet boundary condition

$$\partial_t u = D \,\partial_{xx} u + f(x, u) \,\partial_x u + g(x, u), \qquad u(0) = u(1) = 0$$
 (1.1)

on the interval J = [0, 1]. Here $u = (u_1, \ldots, u_m)$, and f and g are sufficiently regular matrix functions and vector functions, respectively. We assume that the constant coefficient matrix D is similar to a diagonal matrix with positive eigenvalues. In the case of $D = \text{diag}\{d_1, \ldots, d_m\}$ with $d_j > 0$, we deal with reaction-diffusion-convection equations. Under appropriate conditions on f and g, system (1.1) induces a smooth dissipative semiflow $\{\Phi_t\}_{t\geq 0}$ in the phase space $X^{\alpha} \subset C^1(J, \mathbb{R}^m)$ with an appropriate $\alpha > 0$, where $\{X^{\alpha}\}_{\alpha\geq 0}$ is the Hilbert half-scale [1] generated by the linear sectorial operator $u \to -Du_{xx}$ in $X = L^2(J, \mathbb{R}^m)$. In this situation, there exists a global attractor [2]–[4] (simply an attractor below), i.e., a connected compact invariant set $\mathscr{A} \subset X^{\alpha}$ of finite Hausdorff dimension which uniformly attracts bounded subsets X^{α} as $t \to +\infty$.

Our goal is to find conditions under which the dynamics on attractor (final dynamics) of the parabolic system (1.1) is finite-dimensional in the sense of [5]. This means that, for some ODE $\partial_t \xi = h(\xi)$ in \mathbb{R}^N with Lipschitz vector field h, resolving flow $\{\Theta_t\}$, and invariant compact set $\mathscr{K} \subset \mathbb{R}^N$, the phase semiflows $\{\Phi_t\}_{t\geq 0}$ on \mathscr{A} and $\{\Theta_t\}_{t\geq 0}$ on \mathscr{K} are Lipschitz conjugate. In this connection, one can speak [6] about a finite reduction of the evolution problem (1.1).

The main result of the paper (Theorem 4.3) ensures that the final phase dynamics of system (1.1) is finite under the *matching condition*

$$Df(x,u) = f(x,u)D,$$
 $(x,u) \in J \times \operatorname{co} \mathscr{A},$ (1.2)

where $\cos \mathscr{A}$ is the convex hull of \mathscr{A} .

It is known [7] that, in the case of scalar diffusion (D = dE with unit matrix E) and sufficiently regular f = f(u) and g = g(u), there exists an inertial manifold (IM) which is a finite-dimensional

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invariant C^1 -surface in the phase space containing an attractor and exponentially attracting (with an asymptotic phase) all trajectories of the system as $t \to +\infty$. The presence of IM implies the finite-dimensionality of the final dynamics; an extensive literature is devoted to the existence of such manifolds (see, e.g., [3], [4], [6], [8]). An original approach to such problems is presented in recent Anikushin's works (see [9] and the references therein).

For the periodic case (J is a circle of length 1), conditions for the finite-dimensionality of the final dynamics of system (1.1) with D = diag were obtained by the author in [10, p. 13409]. We note that, in the class of periodic systems (1.1) with scalar diffusion, the first example of semilinear parabolic equation of mathematical physics that does not exhibit such a dynamics was constructed in [11, Theorem 1.2].

2. PRELIMINARIES

In what follows, if necessary, we will use the technique developed in [10]. All preliminary constructions in Secs. 2, 3 are performed for the case D = diag. We write system (1.1) as a semilinear parabolic equation (SPE)

$$\partial_t u = -Au + F(u) \tag{2.1}$$

in the *real* Hilbert space $X = L^2(J, \mathbb{R}^m)$ with norm $\|\cdot\|$. Here $A: u \to -Du_{xx}$ with the Dirichlet boundary condition and nonlinearity $F: u \to f(x, u) \partial_x u + g(u)$. For a linear positive definite operator A, we set $X^{\alpha} = \mathscr{D}(A^{\alpha})$ with $\alpha \ge 0$ and $X_0 = X$; then $\|u\|_{\alpha} = \|A^{\alpha}u\|$. We will say that a function Fbelongs to the *class* $W^2(X^{\alpha}, X)$ if

$$F \in C^2(X^{\alpha}, X) \cap \operatorname{Lip}(X^{\alpha}, X)$$
 and $||F(u)|| \le M$ for $u \in X^{\alpha}$ (2.2)

for some $\alpha \in [0, 1)$. In this case, SPE (2.1) generates [1] a smooth compact resolving semiflow $\{\Phi_t\}_{t\geq 0}$ in the phase space X^{α} . Assumption (2.2) implies [8, Lemma 1.1] the X^{α} -dissipativity of (2.1):

$$\limsup_{t \to +\infty} \|\Phi_t u\|_{\alpha} \le r$$

for some r > 0 uniformly in $u \in$ balls in X^{α} . Under these conditions, there exists [2]–[4] a compact attractor $\mathscr{A} \subset X^{\alpha}$ consisting of all bounded complete trajectories $\{u(t)\}_{t \in \mathbb{R}} \subset X^{\alpha}$. In fact, $\mathscr{A} \subset X^{1}$ owing to the *smoothing action* of the parabolic equation [1]. Simple reasoning [10, p. 13410] shows that, in all constructions related to SPE (2.1), the nonlinearity exponent α can be replaced by any value $\alpha_{1} \in (\alpha, 1)$, and if condition (2.2) is satisfied in the two spaces $(X^{\theta}, X^{\theta+\alpha})$ with $\theta > 0$ instead of (X, X^{α}) , then all the above-listed properties of the dynamics remain valid for the phase space $X^{\theta+\alpha}$. In what follows, we will use functions $Y_{1} \to Y_{2}$ of class (2.2) for some Banach spaces Y_{1} and Y_{2} .

As in [10], we will use sufficient conditions for the final finite-dimensionality of the dynamics [12]. Let G(u) = F(u) - Au be the vector field (2.1), and let $\mathcal{N} = \mathscr{A} \times \mathscr{A}$ and Y be Banach spaces.

Definition 2.1[12]. A continuous field $\Pi: \mathcal{N} \to Y$ is said to be *regular* if, for any $u, v \in \mathscr{A}$, the function $\Pi(\Phi_t u, \Phi_t v): [0, +\infty) \to Y$ is of class C^1 with the derivative $\partial_t \Pi(u, v)$ at zero uniformly bounded in $(u, v) \in \mathcal{N}$.

The smoothness of the semiflow $\{\Phi_t\}$ and the invariance of the compact set $\mathscr{A} \subset X^{\alpha}$ imply the regularity of the identical embedding $\mathscr{N} \to X^{\alpha} \times X^{\alpha}$, and hence of any field $\Pi \colon \mathscr{N} \to Y$ that can be continued to a C^1 -mapping in the $(X^{\alpha} \times X^{\alpha})$ -neighborhood of the set \mathscr{N} . In this situation, $\partial_t \Pi(u, v) = \Pi'(u, v)(G(u), G(v))$, where $(\cdot)'$ is the Fréchet differentiation. Under condition (2.2) on the nonlinearity F, the function $u \to G(u)$ on \mathscr{A} is continuous and even Hölder [5] in the X^{α} -metric. The regular fields $\mathscr{N} \to Y$ form a linear structure as well as a multiplicative one if Y is a Banach algebra. In the last case, if all elements $\Pi(u, v) \in Y$ are invertible, then the field Π^{-1} also turns out to be regular.

We start from the decomposition

$$G(u) - G(v) = (T_0(u, v) - T(u, v))(u - v), \qquad (u, v) \in \mathcal{N},$$
(2.3)

where $T_0 \in \mathscr{L}(X^{\alpha})$ and $T \in \mathscr{L}(X^1, X)$ are unbounded linear operators in X similar to positive definite ones. By

$$\Sigma_T = \bigcup_{u,v \in \mathscr{A}} \operatorname{spec} T(u,v)$$

we denote the total spectrum of the operators T.

We will need the special case of [12, Theorem 2.8] in the situation $\Sigma_T \subset \mathbb{R}^+$.

Theorem 2.2. Assume that $F \in W^2(X^{\alpha}, X)$ and

$$T(u,v) = S^{-1}(u,v)H(u,v)S(u,v)$$
(2.4)

on \mathcal{N} , where the unbounded self-adjoint linear operators H(u, v) are positive definite in X, the fields $S, S^{-1} \colon \mathcal{N} \to \mathcal{L}(X)$ and $T_0 \colon \mathcal{N} \to \mathcal{L}(X^{\alpha}, X)$ are regular, and the field $T_0 \colon \mathcal{N} \to \mathcal{L}(X^{\alpha})$ is bounded. Moreover, if the set $\mathbb{R}^+ \setminus \Sigma_T$ contains intervals $(a_k - \xi_k, a_k + \xi_k)$ with $a_k > \xi_k > 0$ such that

$$\xi_k \to \infty, \qquad a_k^{\alpha/2} = o(\xi_k)$$

$$\tag{2.5}$$

as $k \to +\infty$, then the final X^{α} -dynamics of SPE (2.1) is finite-dimensional.

We further assume that the matrix function f = f(x, u) and the vector function g = g(x, u) in (1.1) satisfy the following regularity conditions.

Condition (H). The functions f and g of the class C^{∞} on $J \times \mathbb{R}^m$ are compactly supported in u, and f(x,0) = g(x,0) = 0 for x = 0, 1.

By $\mathscr{H}^s = \mathscr{H}^s(J)$ we denote generalized Sobolev L^2 -spaces (spaces of Bessel potentials [1], [13]) of scalar functions on J with arbitrary $s \ge 0$. If s > 1/2, then $\mathscr{H}^s \subset C(J)$ and \mathscr{H}^s is a Banach algebra [13, Sec. 2.8.3]. The differentiation operator acts in the spaces $\partial_x \in \mathscr{L}(\mathscr{H}^{s+1}, \mathscr{H}^s)$. In fact, the X^s are closed subspaces (with equivalent norm) in the spaces $\mathscr{H}^{2s}(J, \mathbb{R}^m)$ of vector-functions, and $X^s = \mathscr{H}^{2s}(J, \mathbb{R}^m)$ for $s \le 1/4$. For s > 1/4, the space X^s consists of elements $u \in \mathscr{H}^{2s}(J, \mathbb{R}^m)$ with u(0) = u(1) = 0.

Now fix an arbitrary $\alpha \in (3/4, 1)$; then $\mathscr{H}^{2\alpha} \hookrightarrow C^1(J)$ and $X^{\alpha} \hookrightarrow C^1(J, \mathbb{R}^m)$, where the symbol \hookrightarrow denotes a linear continuous embedding of function spaces. Let us use necessary embedding theorems [1], [13]. For an arbitrary C^{∞} -function $z: J \times \mathbb{R}^m \to \mathbb{R}$, the mapping $\psi: u \to z(x, u)$ is a function of class W^2 (see (2.2)) from $C^s(J, \mathbb{R}^m)$ to $C^s(J)$ for all $s \in \mathbb{N}$. This implies that $\psi \in W^2(\mathscr{H}^{2\alpha}(J, \mathbb{R}^m), C^1(J))$. Using the embedding $\mathscr{H}^{s+1} \hookrightarrow C^s(J) \hookrightarrow \mathscr{H}^s$, we can conclude that $\psi \in W^2(\mathscr{H}^s(J, \mathbb{R}^m), \mathscr{H}^s(J))$. So $F \in W^2(X^1, X^{1/2})$ for the nonlinear part $F: u \to f(x, u) \partial_x u + g(u)$ of system (1.1). Moreover, $X^{\alpha} \hookrightarrow C^1(J, \mathbb{R}^m) \hookrightarrow C(J, \mathbb{R}^m) \hookrightarrow X$, and hence $F \in W^2(X^{\alpha}, X)$. We also note that $X^{3/2} \hookrightarrow C^2(J, \mathbb{R}^m)$.

We take X^{α} as the phase space of system (1.1). Following [7], we can show that the phase dynamics of (1.1) in X^{α} is dissipative and there exists a global attractor $\mathscr{A} \subset X^{\alpha}$. Since $F \in W^2(X^1, X^{1/2})$, system (1.1) also generates a smooth dissipative phase semiflow in the space X^1 , and the attractor \mathscr{A} is compact in $X^{3/2}$. As above, we denote $\mathscr{N} = \mathscr{A} \times \mathscr{A}$.

Remark 2.3. The phase dynamics of system (1.1) has the following property: if *Y* is a Banach space, then each vector field $\Pi: \mathcal{N} \to Y$ continuous in the $(X^{\alpha} \times X^{\alpha})$ -metric and extendable to a C^1 -mapping $X^1 \times X^1 \to Y$ is regular in the sense of Definition 2.1.

Indeed, the smoothness of a semiflow in X^1 means the smoothness of the mapping

$$(t, u) \to \Phi_t u \colon (0, +\infty) \times X^1 \to X^1.$$

This ensures the regularity of the identity mapping $\mathscr{N} \to X^1 \times X^1$ and hence the regularity of the field Π on \mathscr{N} .

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3. DECOMPOSITION OF THE VECTOR FIELD ON THE ATTRACTOR

We will apply Theorem 2.2 to SPE (1.1) with D = diag and the phase spaces $X^{\alpha}, \alpha \in (3/4, 1)$. By \mathbb{M}^m we denote the algebra of numerical $(m \times m)$ -matrices with Euclidean norm, and by $Y(J, \mathbb{M}^m)$, the linear spaces of such matrices with elements from some Banach space Y of scalar functions on J = [0, 1]. Following [10, pp. 13412–13413], we set

$$B_0(x;u,v) = \int_0^1 (f_u(x,w(x))w_x(x) + g_u(x,w(x)) d\tau, \qquad (3.1)$$

$$B(x; u, v) = \int_0^1 f(x, w(x)) d\tau$$
(3.2)

for $u, v \in X^{\alpha}$, where $w(x) = \tau u(x) + (1 - \tau)v(x)$, $x \in J$. The entries of the matrices B_0 and B are continuous functions, and for $u, v \in \mathscr{A}$ they are functions of the class C^2 on J. Using the C^1 -smoothness of the mappings $(u, v) \to f_u(x, w)w_x + g_u(x, w)$, $(u, v) \to f(x, w)$, $X^{\alpha} \times X^{\alpha} \to C(J, \mathbb{M}^m)$ for a fixed $\tau \in [0, 1]$ and differentiating the expressions for B_0 and B with respect to the parameter (u, v), we conclude that the mappings

$$(u,v) \to B_0(\cdot; u, v), \qquad (u,v) \to B(\cdot; u, v)$$

$$(3.3)$$

are of the class $C^1(X^{\alpha} \times X^{\alpha}, C(J, \mathbb{M}^m))$. We use the integral mean value theorem for nonlinear operators to write the decomposition of the vector field (1.1) on the attractor $\mathscr{A} \subset X^{\alpha}$ in the form

$$G(u) - G(v) = -Ah + \left(\int_0^1 F'(\tau u + (1 - \tau)v) d\tau\right)h$$

= $Dh_{xx} + B_0(x; u, v)h + B(x; u, v)h_x, \quad u, v \in \mathscr{A},$

where h = u - v, $\tau u + (1 - \tau)v \in co \mathscr{A}$, and $(\cdot)'$ is the Frechet differentiation. To eliminate the dependence on h_x , we (following [14]) apply the transformation $h = U\eta$, where the $(m \times m)$ -matrix function $U(x) = U(x; u, v), x \in [0, 1]$, is the solution of the Cauchy problem

$$U_x = -\frac{1}{2}D^{-1}B(x)U, \qquad U(0) = E.$$
(3.4)

As a result, we obtain relation (2.3) with linear operators

$$T_0(u,v)h = \left(B_0(x) - \frac{1}{2}B_x(x) - \frac{1}{4}B(x)D^{-1}B(x)\right)h,$$
(3.5)

$$T(u,v)h = -DU\,\partial_{xx}U^{-1}h. \tag{3.6}$$

Note that, under the change of variable $h = U\eta$, the Dirichlet boundary conditions for the linear part of (1.1) are preserved. When writing the matrices B_0 , B, U, U^{-1} , we often omit the dependence on u, v, and sometimes, on x.

Lemma 3.1. The field of the operator T_0 on \mathcal{N} is regular with values in $\mathscr{L}(X^{\alpha}, X)$ and bounded with values in $\mathscr{L}(X^{\alpha})$.

Proof. We set $T_0h = Q(x; u, v)h$ in (3.5) with $h \in \mathscr{A} - \mathscr{A} \subset X^{\alpha}$. The convex hull of the attractor \mathscr{A} is bounded in the $X^{3/2}$ -norm equivalent to the norm in $\mathscr{H}^3(J, \mathbb{R}^m)$, and therefore, the matrix functions B, $BD^{-1}B$, and B_0 are bounded uniformly in $(u, v) \in \mathscr{N}$ in $\mathscr{H}^3(J, \mathbb{M}^m)$ and $\mathscr{H}^2(J, \mathbb{M}^m)$, respectively. Thus, the matrix functions B_x and Q are bounded on \mathscr{N} in the norm of $\mathscr{H}^2(J, \mathbb{M}^m)$, and T_0 is the operator of multiplication of vector functions in $X^{\alpha} \subset \mathscr{H}^{2\alpha}(J, \mathbb{R}^m)$ by the matrix $Q \in \mathscr{H}^{2\alpha}(J, \mathbb{M}^m)$ with $2\alpha \in (3/2, 2)$. Since $\mathscr{H}^{2\alpha}(J)$ is a Banach algebra, we see that $T_0(u, v) \in \mathscr{L}(X^{\alpha})$ and $\|T_0(u, v)\|_{\alpha} \leq \text{const on } \mathscr{N}$.

In view of Remark 2.3 and the above-noted smoothness of the mapping (3.3), the regularity of the field of the operator $T_0: \mathscr{N} \to \mathscr{L}(X^{\alpha}, X)$ can be established exactly just as for the case of periodic boundary conditions in [10, Lemma 3.3].

The matrix function U(x) in the Cauchy problem (3.4) can be treated as a bounded linear operator in X.

Lemma 3.2. The fields of the operators $U, U^{-1}: \mathcal{N} \to \mathcal{L}(X)$ are regular.

Proof. For the field U, this can proved as the similar assertion in the periodic case [10, Lemma 3.4]. At the same time, the regularity of U implies the regularity of the field of inverse operators U^{-1} .

Now let $d_{-} = \min_{1 \le j \le m} d_j$ and $d_{+} = \max_{1 \le j \le m} d_j$ for $D = \text{diag}\{d_j\}$. Also let $\{\lambda_n : \lambda_1 < \lambda_2 < \cdots\}$ be the eigenvectors of the linear operator $A = -D \partial_{xx}$. Since

$$\operatorname{spec}(A) = \{ d_j \pi^2 \nu^2, \, \nu \in \mathbb{N}, \, j \in 1, \dots, m \},$$
(3.7)

we have $\lambda_n \leq \pi^2 d_+ n^2$. Using the counting function for spec(A), we obtain

$$n \leq \sum_{j=1}^{m} \frac{\sqrt{\lambda_n}}{\pi \sqrt{d_j}} \leq \frac{m}{\pi \sqrt{d_-}} \sqrt{\lambda_n},$$

and hence

$$\frac{\pi^2 d_-}{m^2} n^2 \le \lambda_n \le \pi^2 d_+ n^2, \qquad n \in \mathbb{N}.$$
(3.8)

Lemma 3.3. *The following estimate holds:*

$$\limsup_{n \to \infty} n^{-1} (\lambda_{n+1} - \lambda_n) > 0.$$

Proof. If, on the contrary, $\lambda_{n+1} - \lambda_n = \beta_n n$ with $\beta_n \xrightarrow{n \to \infty} 0$, then

$$n^{-2}\lambda_n = n^{-2} \left(\lambda_1 + \sum_{k=1}^{n-1} (\lambda_{k+1} - \lambda_k) = n^{-2} \left(\lambda_1 + \sum_{k=1}^{n-1} \beta_k k \right) \\ \leq n^{-2} \left(\lambda_1 + \sum_{k=1}^{n-1} \beta_k n \right) \leq n^{-2}\lambda_1 + n^{-1} \sum_{k=1}^n \beta_k.$$

But this implies the relation $\lambda_n = o(n^2)$, which contradicts the left inequality in (3.8).

4. MAIN RESULTS

By the assumptions of Theorem 2.2, it is necessary, for the operators T(u, v) in (3.6), to establish "uniform" similarity of the positive definite form (2.4) as well as the required sparsity (2.5) of their total spectrum Σ_T . We assume that the regularity Condition (H) is satisfied for the functions f and g in (1.1).

Theorem 4.1. If the matrix D has the form $D = \text{diag}\{d_j\}$ with $d_j > 0$ and the matching condition (1.2) is satisfied, then the phase dynamics on the attractor is finite-dimensional.

Proof. The operator $A = -D \partial_{xx}$ with the Dirichlet condition is self-adjoint and positive definite in X. Assumption (1.2) implies (for any $x \in J$ and $u, v \in \mathscr{A}$) the relation DB(x) = B(x)D for the matrices B(x) = B(x; u, v) in (3.2). Thus, the matrices B(x) and $D^{-1}B(x)$ inherit the block (with respect to the same d_j) structure of the diffusion matrix $D = \text{diag}\{d_1, \ldots, d_m\}$. Therefore, this is also true for the solutions U(x) of the Cauchy problem (3.4), and hence $DU(x) = U(x)D, x \in J$, and

$$T(u,v) = U(u,v)(-D\,\partial_{xx})U^{-1}(u,v)$$

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in (3.6). Thus, for T(u, v) we have the representation (2.4) with $S(u, v) = U^{-1}(u, v)$ and $H(u, v) \equiv A$. The total spectrum Σ_T coincides with spec(A) in (3.7). By Lemma 3.3, there exists an $\varepsilon > 0$ and an increasing sequence of indices n(k) such that $\lambda_{n(k)+1} - \lambda_{n(k)} > \varepsilon n(k)$ for $k \ge k_0$. Set

$$a_k = \frac{\lambda_{n(k)+1} + \lambda_{n(k)}}{2}, \qquad \xi_k = \frac{\lambda_{n(k)+1} - \lambda_{n(k)}}{3}, \qquad M = \pi^2 d_+$$

From the right inequality in (3.8), we obtain

$$a_k \le M\left(n^2(k) + n(k) + \frac{1}{2}\right) \le 3Mn^2(k) \le \frac{3M}{\varepsilon^2} (\lambda_{n(k)+1} - \lambda_{n(k)})^2 \le \frac{27M}{\varepsilon^2} \xi_k^2$$

for $k \ge k_0$; i.e., $a_k = O(\xi_k^2)$ as $k \to \infty$. Since $a_k^{\alpha/2} = o(\xi_k)$ for $\alpha \in (3/4, 1)$ as $k \to \infty$, we see that the desired assertion follows from Lemmas 3.1 and 3.2 and Theorem 2.2.

Remark 4.2. Parabolic systems (1.1) with D = diag demonstrate a finite-dimensional dynamics on the attractor for any admissible nonlinearities f and g in the case of scalar diffusion and under the condition f = diag in the case of m distinct diffusion coefficients d_j . In the case of s distinct diffusion coefficients with 1 < s < m, the dynamics on the attractor is finite-dimensional under the condition that the matrix function f inherits the block structure (with the same d_j) of the matrix $D = \text{diag}\{d_j\}$.

Now we state the main result. We assume that the matrix D in system (1.1) has the form $D = C\overline{D}C^{-1}$, where the matrix C is nonsingular and $\overline{D} = \text{diag}\{d_1, \ldots, d_m\}$ with $d_j > 0$. The linear operator $-D \partial_{xx} = -C(\overline{D} \partial_{xx})C^{-1}$ is sectorial in $X = L^2(J, \mathbb{R}^m)$. The linear change of variables u = Cv reduces (1.1) to the system of equations

$$\partial_t v = \overline{D} \,\partial_{xx} v + \overline{f}(x, v) \partial_x v + \overline{g}(x, v), \qquad v(0) = v(1) = 0,$$

$$\overline{f}(x, v) = C^{-1} f(x, Cv)C, \qquad \overline{g}(x, v) = C^{-1} g(x, Cv).$$
(4.1)

The matrix function \overline{f} and the vector function \overline{g} inherit the regularity properties (H) of the original functions f and g. The phase semiflows of systems (4.1) and (1.1) are linearly conjugate. System (4.1) is dissipative in X^{α} , and hence this is also true for system (1.1). The attractors \mathscr{A} of system (1.1) and $\overline{\mathscr{A}}$ of system (4.1) are related by the expression $\mathscr{A} = C\overline{\mathscr{A}}$. According to the definition of the finite-dimensionality of the phase dynamics (Sec. 1), systems (4.1) and (1.1) demonstrate this property simultaneously.

Theorem 4.3 (main). If the matrix D is similar to diag $\{d_j\}$ with $d_j > 0$ and the matching condition (1.2) is satisfied, then the phase dynamics of system (1.1) is finite-dimensional.

Proof. Since Df(x, u) = f(x, u)D on $J \times co \mathscr{A}$, it follows that

$$\overline{D}\overline{f}(x,v) = C^{-1}DC \cdot C^{-1}f(x,Cv)C = C^{-1}Df(x,u)C$$
$$= C^{-1}f(x,u)DC = C^{-1}f(x,Cv)C\overline{D} = \overline{f}(x,v)\overline{D}$$

on $J \times \operatorname{co} \overline{\mathscr{A}}$. Here $u \in \operatorname{co} \mathscr{A}$ and $v \in \overline{\mathscr{A}}$. As we can see, condition (1.2) is satisfied for the matrix function \overline{f} , and the dynamics of system (4.1) on the attractor $\overline{\mathscr{A}} \subset X^{\alpha}$ is finite-dimensional by Theorem 4.1. This also implies the finite-dimensionality of the dynamics of system (1.1) on the attractor $\mathscr{A} \subset X^{\alpha}$.

Remark 4.4. Under the matching condition (1.2), the final dynamics of system (1.1) is finite-dimensional if all eigenvalues of the matrix D are distinct and positive. Condition (1.2) is satisfied, in particular, for $f = D_1 \varphi$, where the numerical matrix D_1 commutes with D and $\varphi = \varphi(x, u)$ is a smooth scalar function compactly supported in u.

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