

12th International Conference on Magnetic Fluids

## A magnetic fluid bridge between coaxial cylinders with a line conductor

V.A. Naletova<sup>a,b</sup>, V.A. Turkov<sup>b</sup>, A.S. Vinogradova<sup>a,b,\*</sup>

<sup>a</sup>Department of Mechanics and Mathematics, Lomonosov Moscow State University, Moscow, 119992, Russia

<sup>b</sup>Institute of Mechanics, Lomonosov Moscow State University, Michurinsky Pr. 1, Moscow, 119192, Russia

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### Abstract

The behavior (break-up and creation) of a magnetic fluid bridge between two coaxial cylinders in a magnetic field of a line conductor for any values of wetting angles of a magnetic fluid is investigated.

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*Keywords:* Magnetic fluid; Line conductor

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### 1. Introduction

A surface shape of a magnetic fluid (MF) changes spasmodically in a magnetic field of a line conductor [1, 2, 3]. In [1] a MF drop on a line conductor was considered. In [2, 3] the behavior of a MF between two coaxial cylinders (there is a line conductor in the bulk of the inner cylinder) was investigated when MF wets cylinders (wetting angles are no more than  $\pi/2$ ). In the present paper the behavior of a MF bridge between two coaxial cylinders is investigated for any values of wetting angles.

### 2. Surface shape of MF near a line conductor

Consider a MF bridge ( $V_m$  is the volume of a MF) between two coaxial cylinders with circular cross section and the radii  $R_c$  and  $r_0$ ,  $R_c > r_0$ , Fig.1. A line conductor with current  $I$  is along the axis of cylinders. Let the MF be a heavy, incompressible fluid. The MF is immersed in a liquid with the same density. The surface of a MF is  $z = h(r)$ ,  $r^2 = x^2 + y^2$  and the magnetic field of a conductor  $H = 2I / (cr)$  is not deformed. The magnetization of a MF is described by the Langevin's law:  $M(\xi) = M_S L(\xi)$ ,  $L(\xi) = cth(\xi) - 1/\xi$ ,  $\xi = mH / (kT)$ ,  $m = M_S/n$ . The surface of a MF near the conductor in the non-dimensional form is [3] ( $P^*(r^*) = \ln(sh(\xi_0 H^*)/\xi_0 H^*)$ )

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\* Corresponding author. Tel./fax: +7-495-939-5974.

E-mail address: [vinogradova-as@mail.ru](mailto:vinogradova-as@mail.ru)

$$h^*(r^*) = \int \frac{G}{(1-G^2)^{1/2}} dr^*, \tag{1}$$

$$G(r^*) = \frac{C}{r^*} + Br^* - \frac{P_1}{r^*} \int_1^{r^*} r^* P^* dr^*.$$

Here  $r^* = r/r_0$ ,  $h^* = h/r_0$ ,  $H^* = H/H_0$ ,  $H_0 = 2I/(cr_0)$ ,  $\xi_0 = mH_0/(kT) = 2mI/(cr_0kT)$ ,  $P_1 = nkTr_0/\sigma$ ,  $R_c^* = R_c/r_0$ . Later the signs “\*” are omitted and parameters are considered as non-dimensional. If the MF bridge exists the values of constants  $B$  and  $C$  are determined from boundary conditions:  $G(r = 1) = -\cos(\theta_1)$ ,  $G(r = R_c) = -\cos(\theta_2)$  (here  $\theta_1, \theta_2$  are wetting angles)  $C = -\cos(\theta_1) - B$ ,

$$B = (P_1 \int_1^{R_c} rP dr + \cos(\theta_1) - \cos(\theta_2)R_c)/(R_c^2 - 1).$$

In Fig.1 the MF is represented in case of non-wetting a) before break-up and b) after break-up.

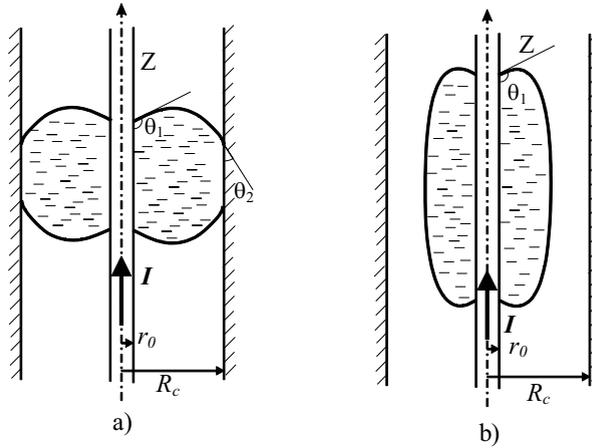


Figure 1: The MF in case of non-wetting a) before break-up and b) after break-up

### 3. Break-up of the MF bridge between two cylinders

A critical value of the Langevin’s parameter  $\xi_c$  exists when  $G = -1$ ,  $dG/dr = 0$ . For  $\xi_0 > \xi_c$  the MF bridge between cylinders can not exist. When  $\cos \theta_2 = -\cos \theta_1$ , the value  $\xi_c$  depends on  $P_1, R_c, \cos \theta_1$ . The calculation  $\xi_c = \xi_c(\cos \theta_1)$  and other calculations were done numerically for  $P_1 = 5.65$  ( $r_0 = 0.1$  cm,  $T = 300^\circ$  K,  $n = 0.95 \cdot 10^{17}$  cm<sup>-3</sup>,  $\sigma = 70$  din/cm) and different values of  $R_c$  (Fig.2). It may be seen that  $\xi_c$  depends on  $\cos \theta_1$  non-monotonically and the single maximum of  $\xi_c$  exists for  $\theta_1 > \pi/2$  (the case of non-wetting). For  $\xi_0 = \xi_{ob}$ ,  $\xi_{ob} < \xi_c$  the bridge break-up occurs in the moment of the touch of upper and lower surfaces at some point  $r = r_1$  and when the MF volume  $V_m$  is equal to a minimal MF volume between cylinders  $V_0$ ,  $V_m = V_0(\xi_{ob})$  (this equality is used for the determination of  $\xi_{ob}$ ). The dependences  $V_0 = V_0(\xi_0)$  in case of non-wetting are monotonic (Fig.3). The dependence  $V_0 = V_0(\xi_0)$  in case of wetting is non-monotonic (Fig.4). It means that in case of wetting the bridge break-up and the formation of MF volumes  $V_1$  and  $V_2$  on inner and outer cylinders may occur for both the increase and the decrease of current, when  $V_m < V_0(\theta)$ . In case of non-wetting the bridge break-up (the formation of a drop on the conductor) may occur only if the current increases.

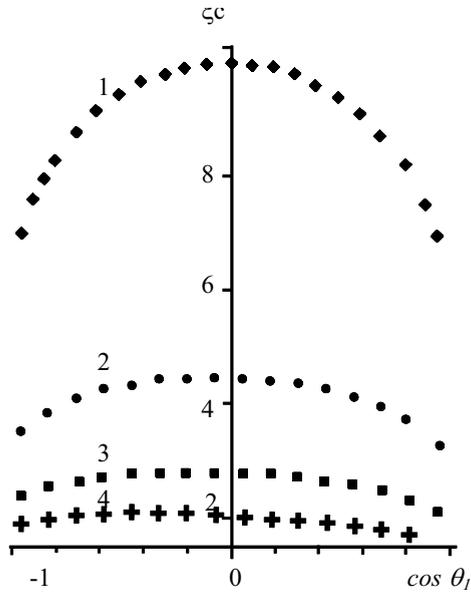


Figure 2: The value of  $\xi_c$  versus  $\cos \theta_l$  for 1 –  $R_c=1.5$ , 2 –  $R_c=2$ , 3 –  $R_c=3$ , 4 –  $R_c=6$

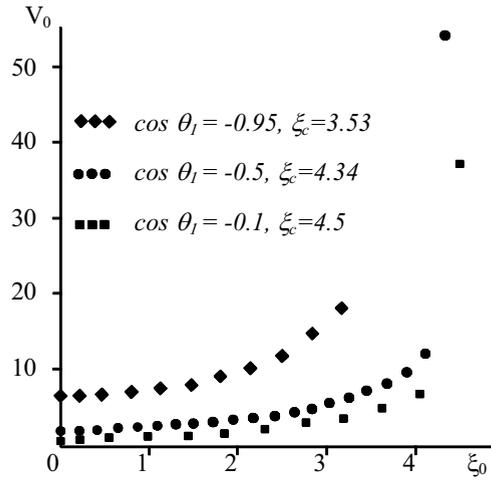


Figure 3: The value of  $V_0$  versus  $\xi_0$  for  $\cos \theta_l < 0$ ,  $R_c=2$

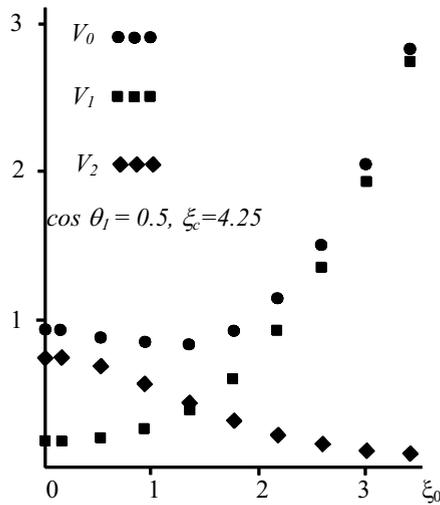


Figure 4: Values of  $V_0$ ,  $V_1$ ,  $V_2$  versus  $\xi_0$  for  $\cos \theta_l > 0$ ,  $R_c = 2$

#### 4. MF drop on a line conductor

After the bridge break-up for some  $\xi_{0b}$ ,  $\xi_{0b} < \xi_c$  a MF drop appears on a line conductor with the current  $I_b$  of the radius  $r_0$  (Fig.1 b). In case of wetting the drop volume  $V_d$  is equal to  $V_1(\xi_{0b})$  and in case of non-wetting the drop volume  $V_d$  is equal to  $V_m$ . The drop shape is described by the formula (1). The constant  $C$  is determined as  $C = -\cos \theta_l - B$  from the condition  $G(r = l) = -\cos \theta_l$ . The dependence  $B = B(r_d)$  ( $r_d$  is the thickness of the drop) may be obtained from the equation  $G(r_d) = -1$ . The thickness of the drop  $r_d$  is not given. But we know the value of the drop volume  $V_d$ . The dependences  $V = V(r_d)$  for different  $\xi_0$  (here  $V$  is the volume of some drop) in case of non-wetting ( $\cos \theta_l = -0.5$ ) were done numerically and are shown in Fig.5. It is shown that there are three critical values of  $\xi_0$ :  $\xi_{01}$ ,  $\xi_{02}$ ,  $\xi_{03}$ , when the dependence  $V(r_d)$  strongly changes. If  $\xi_0 < \xi_{01}$  the dependence  $V = V(r_d)$  is a monotonous one (lines 1, 2 in Fig.5). If  $\xi_{01} < \xi_0 < \xi_{02}$  the dependence has a maximum and a minimum, and some values of the volume  $V$  correspond to three values of the drop thickness  $r_d$  (lines 3, 4 in Fig.5). If  $\xi_{02} < \xi_0 < \xi_{03}$  the dependence has two branches, some values of the volume  $V$  correspond to three values of the drop thickness  $r_d$  (lines 5, 6 in Fig.5). If  $\xi_0 > \xi_{03}$  the dependence  $V = V(r_d)$  is a monotonous one (lines 7, 8 in Fig.5). So for the big enough drops and for  $\xi_0 > \xi_{01}$  spasmodic changes of the drop radius and hysteresis phenomena can be observed for cyclic increase and decrease of current.

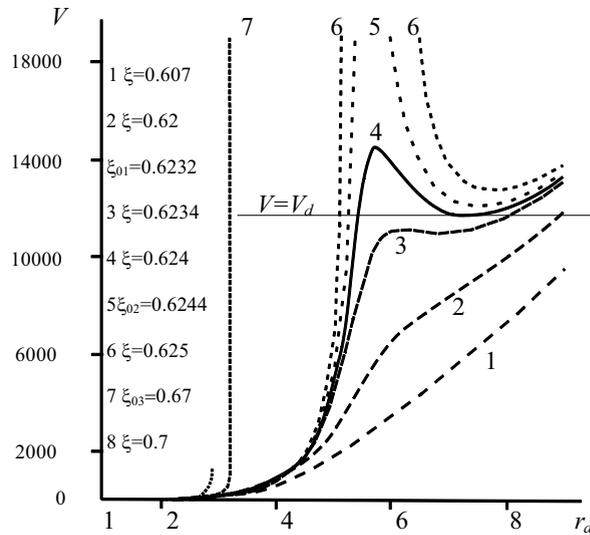


Figure 5: The dependences  $V = V(r_d)$  ( $\cos\theta_1 = -0.5$ ,  $P_1 = 42.3$ ,  $r_0 = 0.065$  cm,  $T = 300^\circ$  K,  $\sigma = 12$  din/cm)

## 5. Formation of the MF bridge between two cylinders in case of non-wetting

Let the current increases from zero to  $I_b$  ( $\xi_0 = \xi_{0b}$ ) for which the break-up of the MF bridge with volume  $V_m$  occurs. In case of non-wetting the MF drop of the volume  $V_m$  appears on a line conductor after the break-up. We can solve the problem about formation of the MF bridge after its break-up when the current decreases from  $I_b$  to zero using the dependence  $V = V(r_d, \xi_0)$ , Fig.5. If the inequality  $R_c > r_d(I=0)$  is valid the MF bridge can not appear. The value of  $r_d(I=0)$  is the solution of the equation:  $V_m = V(r_d, 0)$ . We can find geometrically the value of  $r_d(I=0)$ :  $r_d(I=0)$  is an abscissa of the point of intersection of lines  $V = V_m$  and  $V = V(r_d, 0)$ . The minimal volume  $V_m^{(\min)}$  for which the MF bridge appears is equal to  $V_m^{(\min)} = V(R_c, 0)$ . If the inequality  $R_c > r_d(I=0)$  is not valid the MF bridge between two cylinders can appear for some current  $I_f$ ,  $0 < I_f < I_b$ . The value of  $\xi_{0f} = \xi_0(I_f)$  is the solution of the equation  $V_m = V(R_c, \xi_{0f})$ . This value can be found geometrically:  $\xi_{0f}$  is the Langevin's parameter for which a line  $V = V(r_d, \xi_{0f})$  goes across the point with coordinates  $(R_c, V_m)$  on the plane  $(r_d, V)$ .

## 6. Conclusion

It is shown that a break-up of the MF bridge between cylinders is for  $\xi_{0b} < \xi_c$  when the current increases from zero to  $I_b$ . The volumes of the MF which appear on cylinders,  $\xi_c$  and  $\xi_{0b}$  are calculated numerically. The problem about a MF drop on a line conductor is solved for arbitrary wetting angles. The dependences of the volume of some drop on the thickness of the drop  $V = V(r_d)$  for different  $\xi_0$  are calculated. It is shown that there are three critical values of the parameter  $\xi_0$  when the dependence  $V(r_d)$  strongly changes. The least critical value  $\xi_{01}$  is found for the first time. The dependence  $V = V(r_d)$  is necessary for the solution of the drop problem and for the solution of the problem of formation of the MF bridge between two cylinders after its break-up when the current decreases from  $I_b$  to zero. It is shown that the MF bridge can appear if  $V_m \geq V_m^{(\min)}$  or can not appear if  $V_m < V_m^{(\min)}$ , when the current decreases from  $I_b$  to zero. The minimal volume  $V_m^{(\min)}$  of the MF for which the MF bridge appears is found for different wetting angles.

## Acknowledgements

This work was the part of a study supported by the RFBR (project No. 10-01-90001).

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